# Many-body magic: tensor network states and gauge theories





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References: PRXQ 4, 040317 (2023), PRA 109, L040401 (2024), PRL 133, 010601 (2024), PRB 110, 045101 (2024), 2409.01789 and 2411.11720.

#### Nice to be at this workshop

#### Don't have to justify the name magic

#### Don't have to discuss why that's cool

## Don't have to spend 5 minutes reviewing Clifford group :)

#### Why shall many-body theorists bother about magic?

"Real" cost of quantum computing

Opportunities to learn about quantum phenomena (classification, probing etc.)

Opportunities to harness complexity and leverage this to perform better/new simulations and theory

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Opportunities to learn about quantum phenomena (classification, probing etc.)

Entanglement: topological orders, criticality

Opportunities to harness complexity and leverage this to perform better/new simulations and theory

Entanglement: tensor network simulations

### Few things I'd like to discuss today

Tensor network and magic: how to compute it

Compute

PRXQ 4, 040317 (2023), PRL 133, 010601 (2024)

Magic in lattice gauge theory: is there something interesting?

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Can magic lead us to understand the deeper structure of many-body Hilbert spaces?

New methods?

PRB 110, 045101 (2024) and 2411.11720.

#### Step 0: what and how we compute

Many meaningful witnesses and measures, interesting quantities

Robustness, mana, min-relative entropy of magic, Bell magic

Our target: stabilizer Renyi entropies (Alioscia's talk)

$$M_n(\rho) = \frac{1}{1-n} \log \sum_{P \in \mathcal{P}_N} \frac{|\operatorname{Tr}(\rho P)|^{2n}}{d^N}$$

Leone, Oliviero & Hamma, PRL 128, 050402 (2022)

### **Stabilizer Renyi entropies**

$$M_n(\rho) = \frac{1}{1-n} \log \sum_{P \in \mathcal{P}_N} \frac{|\operatorname{Tr}(\rho P)|^{2n}}{d^N}$$

Formulated as expectation values of string operators



Satisfies properties of measures\*

Probabilistic interpretation - pure gold

Still, requires to measure of measurements  $\propto 4^N$ 

Leone, Oliviero & Hamma, PRL 128, 050402 (2022). See also discussion in Piroli and Haug, Quantum 2023.

#### How can we compute Stabilizer entropies?

$$M_n(\rho) = \frac{1}{1-n} \log \sum_{P \in \mathcal{P}_N} \frac{|\operatorname{Tr}(\rho P)|^{2n}}{d^N}$$

We focus on one-dimensional (+small 2D) systems

Tons of physics: conformal criticality, topology, frustration, dynamics

- Relevance to synthetic matter experiments
- Paradigms for quantum computing blocks
- Solid theory framework: tensor networks

#### Stabilizer entropies and tensor networks

$$M_n(\rho) = \frac{1}{1-n} \log \sum_{P \in \mathcal{P}_N} \frac{|\operatorname{Tr}(\rho P)|^{2n}}{d^N}$$





#### **Exact approaches**

Haug, & Piroli, PRB 107, 035148 (2023); Tarabunga et al., PRL 133, 010601 (2024).

#### Stochastic

Lami & Collura, PRL '23; Piroli and Haug, Quantum '23; Tarabunga et al., PRXQ 4, 040317 (2023)

## Insight: special structure of matrix product states might work!

Basic idea: use a replicated matrix product state (MPS)



Haug, & Piroli, PRB 107, 035148 (2023).

Very elegant, yields exact SRE

While efficient, cost is

 $O(N\chi^{6n})$ 



Demonstrated only for Ising model at  $\chi = 12$ 

MPS: states with finite (area law) entanglement between partitions

#### New insight: stochastic method

In some cases, direct sampling is possible!

Lami & Collura, PRL '23; Piroli and Haug, Quantum '23

For the case of connected partitions and OBC, with MPS

$$O(N\chi^3)$$



#### Pauli MPS



Matrix product state in a replicated Pauli basis - Pauli MPS PRL 133, 010601 (2024)

- No sampling errors
- Exact control of truncation errors
- Modest scaling  $O(\chi^4)$



#### Pauli MPS

Can be used to measure other magic-related quantities (e.g., Bell Magic - Haug and Kim, 2023)



#### Going around: Pauli Markov chains



$$M^{(n)}(|\psi_N\rangle) = \frac{1}{1-n} \log\left\{\sum_{P \in P_N} \frac{Tr(\rho P)^{2n}}{2^N}\right\}$$

Sample not states, but the **distribution of Pauli strings**, with importance sampling!

#### **Pauli-Markov chains**

NB: straightforwardly applicable to experiments, albeit role of statistical errors unclear. Also applicable to other numerical methods

## **Our tool: Pauli Markov chains**

#### Algorithm:

**Input**: a quantum state  $\rho$  and number of sampling  $N_S$ 

- 1: Initialize the Pauli string P.
- 2: Compute  $Tr(\rho P)$  and  $\Pi_P$ .
- 3: for  $(i = 1; i \le N_S; i + +)$  do
- 4: Propose a candidate Pauli string P'.
- 5: Compute  $\operatorname{Tr}(\rho P')$  and  $\Pi_{P'}$ .
- 6: Accept the move with probability:  $\min\left(1, \frac{\Pi_{P'}}{\Pi_P}\right)$ .
- 7: Measure the estimators.

8: end for

**Output**: a Markov chain of P with probability  $\Pi_P$ .

Key elements to address:

Reliability of

estimators

- Tensor contractions
- Memory effects

#### **Tensor contractions: example of trees**



For local (e.g., 1- or 2sites) updates, can be contracted very efficiently!!!

$$\langle O_1 \cdots O_i \cdots O_N \rangle =$$

 $O(\log(N)\chi^4)$ 

- arbitrary partitions
- PBC
- dimension plays very little role
- Stochastic
- Easily extended to MPS, PEPS etc.



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#### **Memory effects: Autocorrelations**



2D Ising LGT at h=3, N=LxL  $\chi=60, N_S=10^6$ 

#### Method development: summary





We can now do trustworthy computations of magic for relatively large volumes and entangled states

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We can now do trustworthy computations of magic for relatively large volumes and entangled states

Now, we can look at some physics!! *This was not possible before 2022* - a great example of quantum info and many-body synergy!

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### Why the hell gauge theories???

**Fundamental interactions** 

**Topological matter** 

Actually, epitome models for error correction!!



U



## First studies in LGTs: upper bounds

## In HEP, people are very interested in complexity of gauge theory simulations

							Schw	vinger bosons
x	$\eta$	L	$t/a_s$	$\Delta$	$\alpha_{\mathrm{Trot.}}$	$\alpha_{\rm Newt.}$	Qubits	T gates
1	4	100	1	0.01	90%	9%	2626	$8.19713  imes 10^{11}$
1	4	100	1	0.001	90%	9%	2704	$3.09951  imes 10^{12}$
1	4	100	10	0.01	90%	9%	2704	$3.0993\times10^{13}$
1	4	100	10	0.001	90%	9%	2808	$1.2146\times10^{14}$
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#### 2D SU(2), state preparation

Davoudi, Shaw, Stryker, Quantum 2024

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1	4	1000	10	0.01	90%	9%	19008	$1.225$ 54 $\times$ 10 <sup>15</sup>
1	4	1000	10	0.001	90%	9%	19086	$4.4867 \times 10^{15}$

#### 2D SU(2), state preparation

Davoudi, Shaw, Stryker, Quantum 2024

$$H_{Z_2-\text{Gauge}} = -h\sum_{i}\prod_{i\in i}\tau_i^x - \sum_i\tau_i^z$$



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 very different from 1D: magic displays crossing at criticality

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very impressive collapse scaling!

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 very different from 1D: magic displays crossing at criticality

very impressive collapse scaling!





#### Lattice gauge theories



#### Lattice gauge theories



 at finite bond dimension, magic detects critical behavior *better* than the order parameter



#### Now, theory with matter - Schwinger model

#### Quantum link formulation

Chandrasekharan and Wiese, 1994-97; Banerjee et al., PRL 2012

$$H = \sum_{j} \left[ -t(\psi_{j}^{\dagger} S_{j,j+1}^{+} \psi_{j} + \text{h.c.}) + V n_{j} n_{j+1} + U(-1)^{j} n_{j} \right]$$

Exact dual to a constrained spin model ("PXP")

Surace et al., PRX 2020

$${\hat H}_{ ext{FSS}} \!= \sum_{j} ig(\Omega\, {\hat \sigma}_{j}^{x} + \delta\, {\hat n}_{j}ig) + \sum_{j < l} V_{j,l}\, {\hat n}_{j} {\hat n}_{l}$$

Fendley, Sengupta, Sachdev, PRB2004; Phase diagram: Chepiga and Mila, Giudici et al., and more Experiments: Lukin's group

#### 1D gauge theories and constrained spin models



See also: Smith et al., 2406.14348

## 1D gauge theories and constrained spin models



Dashed: integrable lines. Thin: exactly soluble L=30. Error  $< 10^{-3}$ 

See also: Smith et al., 2406.14348

## 1D gauge theories and constrained spin models



L=30. Error < 10^-3

Maximum typically inside phases Minimum inside ordered phases

Dashed: integrable lines. Thin: exactly soluble

See also: Smith et al., 2406.14348

#### Magic approaching the continuum limit





### Magic approaching the continuum limit



Magic shows flexes at transition points



## Magic approaching the continuum limit



- Magic shows flexes at transition points
- Criticality unrelated to magic strength



#### What have we learned?

Q: Does magic relate to physical phenomena?

Yes! Both in 1D and 2D, full state magic diagnoses transition points (flexes), and displays universality. But importantly, one needs to look at derivatives!



This is already quite informative, even more than entanglement!

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<u>Idea</u>: try to understand the relation between magic and separability in the context of simple states



Spin-1 
$$H = \sum_{i=1}^{N} (S_i^x S_{i+1}^x + S_i^y S_{i+1}^y + J_Z S_i^z S_{i+1}^z) + D \sum_{i=1}^{N} S_i^{z^2}$$



<u>Idea</u>: try to understand the relation between magic and separability in the context of simple states





<u>Idea</u>: try to understand the relation between magic and separability in the context of simple states





 In gapped phases, convergence is too fast to understand scaling
 At critical points, we always observe:

$$m_1(\chi) = m_1 + \frac{a_o}{\chi^2}$$

Same scaling also true for random MPS (Lami et al., arXiv:2404.18751)



Clifford augmented MPS! arXiv:2405.09217



### CAMPS: a novel class of variational states

Clifford augmented MPS! arXiv:2405.09217

See also: stabilizer TN (2403.08724) and hybrid stabilizer MPOs (PRL 133, 150604 (2024))

 Fundamentally simple structure (local contractions)

Can input in arbitrary entanglement at fixed cost

 Impressive performance gain for ground states!
 Less clear for dynamics: 2407.01692; 2407.03202



#### CAMPS: a novel class of variational states?

#### Trades a constant overhead... for what?

Is there a new corner of the Hilbert space?
 If yes, when are CAMPS overpowering MPS?

#### Reverse questions: can we disentangle MPS with local Clifford??

### Locally Clifford disentangler

#### Local Stabilizer disentangler: CAMPS cost is constant

Distematically removes entanglement, magic is invariant Stabilizer disentangling protocol



Key methodological advance: number of gates is 20, not 720 as in 2405.09217

#### Is there a new corner of the Hilbert space?

Clifford disentangled state reduced density matrix

$$\rho_A^{(SM)} = \text{Tr}_{\bar{A}} |\Psi[C_{\ell;L-\ell}]\rangle \langle \Psi[C_{\ell;L-\ell}] |$$

Its entropy (denoted as SMEE)

$$S_A^{(SM)} = -\mathrm{Tr}\rho_A^{(SM)} \ln \rho_A^{(SM)}$$

Entanglement gain:

$$\Delta(\ell_A; L) = S_A - S_A^{SM}$$

#### Is there a new corner of the Hilbert space?

Clifford disentangled state reduced density matrix

$$\rho_{A}^{(SM)} = \operatorname{Tr}_{\bar{A}} |\Psi[C_{\ell;L-\ell}]\rangle \langle \Psi[C_{\ell;L-\ell}]$$
Its entropy (
Does Delta increase with size, or not?
$$S_{A}^{\star} = -\operatorname{Tr}\rho_{A}^{\star} + \operatorname{In}\rho_{A}^{\star}$$

Entanglement gain:

$$\Delta(\ell_A; L) = S_A - S_A^{SM}$$

#### XXZ model



#### XXZ model



XXZ model

**Tricritical Ising chain** 





**Tricritical Ising chain** 



### CAMPS: a novel class of variational states?

#### 1. Is there a new corner of the Hilbert space?



Local Clifford disentanglable (LCD)

$$\chi_{MPS} - \chi_{CAMPS} = O(L^{\beta})$$

2. If yes, when are CAMPS overpowering MPS?

### When are CAMPS better?

Qualitative insight: CAMPS can decode away entanglement stored in non-local stabilizer degrees of freedom



Extreme case: only stabilizer has support in A+B

—> I shall be able to remove some entanglement between A and B by keeping magic constant

—> quantity to look  $L_n(\rho_{AB}) = \tilde{M}_n(\rho_{AB}) - \tilde{M}_n(\rho_B) - \tilde{M}_n(\rho_A)$ at: mutual SRE

#### What is mutual SRE monitoring

$$L_n(\rho_{AB}) = \tilde{M}_n(\rho_{AB}) - \tilde{M}_n(\rho_B) - \tilde{M}_n(\rho_A)$$

Not a measure of magic in general (at least, not for qubits)

Still, very important! Participation entropy in stabilizer space (Turkeshi, Schiro',...)

If **mutual SRE<0**, there could be some shared qubit that can be disentangled via Clifford only!

#### Can mSRE distinguish?

XXZ model

#### **Tricritical Ising chain**



### CAMPS: a novel class of variational states?

1. Is there a new corner of the Hilbert space?



Local Clifford disentanglable (LCD)

$$\chi_{MPS} - \chi_{CAMPS} = O(L^{\beta})$$

2. If yes, when are CAMPS overpowering MPS?

CAMPS systematically better if mutual SRE becomes negative!! Non-local stabilizer (anti-) correlations

#### **Analytical understanding: Cluster Ising models**

$$H_1 = -\sum_{i=1}^{L-2} (S_i^x S_{i+1}^z S_{i+2}^x) + h \sum_{i=1}^{L} S_i^z$$

c=1 critical point feature an exact self-duality

$$H = \sum_{s=1}^{2} H_{\text{Ising}}^{s}$$

The duality transformation is a Clifford transformation
 Of course, mutual SRE <0! (Hidden Stabilizer info encoded)</li>

Smacchia, Pascazio, Fazio, Amico, .. PRA 2011

#### **Analytical understanding: Cluster Ising models**

**Predictions:** 

(A) At I=L/2, SMEE vanishes

(B) For I<L/2, SMEE coincide with that of a Clifford disentangled Ising chain (central charge from 1 to 1/2)



### **Open questions**

# We are barely scratching the surface of the interface magic/many-body systems!

Simply, too many open questions to list, here some:

Magic and deeper entanglement structures? (Modular Hamiltonian complexity)

Which 2D models can CAMPS solve that MPS cannot?

What about many-body with true measures (robustness)??

More on gauge theories? Dynamics? (See talk by Martin)





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